

The Wikipedia article on weak measurements existing at this time (Nov. 24, 2010) seems to me highly misleading and in some ways simply wrong. I posted a criticism of the article on the “Discussion” page a few days ago, but I am hesitant to attempt to rewrite the article for fear of being drawn into a time-consuming controversy. Were I to rewrite it, it would go something like the following.

If anyone thinks it has sufficient merit to form the basis for a Wikipedia article, I have no objection so long as it is not attributed to me unless presented exactly as I wrote it. I have no doubt that it could be improved, but I am also concerned that I might not agree with some well-meaning “improvements”.

Sample encyclopedia article on quantum weak measurement

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“Weak measurement” is a technique to measure the average value of a quantum observable A without appreciably affecting the initial state s of the system being measured. Weak measurements differ from normal (sometimes called “strong” or “von Neumann”) measurements in two ways:

1. If A has discrete spectrum (which we assume for simplicity), a strong measurement when the system is in state s yields an eigenvalue of A ; if the measurement is repeated many times (starting each time with the system in state s) one obtains a sequence of eigenvalues of A which when averaged yield an approximation to $\langle s|A|s\rangle$, the expectation of A in the state s .

By contrast, a *weak measurement* only yields a sequence of numbers which *average* to $\langle s|A|s\rangle$. For example, a strong measurement of the spin of a spin-1/2 particle must yield spin 1/2 or -1/2, but a particular weak measurement could yield spin 100 [1], while a subsequent weak measurement on an identical system might be -128.3. Typically, a single weak measurement gives little information; only the average of a large number of such measurements is meaningful.

2. A strong measurement changes (“projects”) an initial pure state s to an eigenvector of A . (The particular eigenvector obtained cannot be predicted, though its probability is determined.) This substantially changes the state s unless s happened to be close to that eigenvector.

However, a weak measurement does not substantially change the initial state.

Weak measurements are usually implemented by coupling the original system S to be measured with an auxiliary quantum “meter system” M . The meter

system is often visualized as a macroscopic meter with a pointer that moves along a scale, though in practice various microscopic quantum systems are used.

The composite system is mathematically represented as the tensor product of S with M , denoted $S \otimes M$. A “product” state in this tensor product is typically denoted $|s\rangle|m\rangle$, where s is a state of S and m a state of M . States which are not product states are called *entangled* states. An example of a state which can be shown to be entangled is

$$e := \sqrt{p}|s_1\rangle|m_1\rangle + \sqrt{1-p}|s_2\rangle|m_2\rangle,$$

with s_1, s_2 linearly independent states of S with $|s_1| = |s_2| = 1$, m_1, m_2 orthonormal states of M , and $0 < p < 1$ a probability. The state of S corresponding to e is not a “pure” state, but a mixture of pure states, (described by the density matrix $\text{trace}_M P_e$, where trace_M denotes the partial trace over M and P_e the projector to the subspace spanned by e). It is the mixed state which is in pure state s_1 with probability p and in pure state s_2 with probability $1 - p$. If p is close to 0 or 1, then one could say that e is “slightly entangled”, in which case the corresponding state of S is close to s_2 or s_1 , respectively. (If s_1 and s_2 are not independent, then e is not entangled, and the corresponding state of s is s_1 .)

Weak measurements of an observable A when the system S is in pure state s_1 can be implemented by putting the composite system in a state like e with p close to 1 and measuring an observable B in the meter system which has eigenvectors m_1, m_2 . The measurement will be “weak”, in the sense that the state of S after the measurement will differ little from s_1 . By suitably choosing the observable B , and the state s_2 , it can be arranged that the average value of B when the composite system is in the state e equals (to arbitrary accuracy) the average value of A when S is in state s_1 .

To get an idea of how this can work, let ϵ be a small positive parameter and $B = B(\epsilon)$ an observable on M with $Bm_1 = \frac{1}{2\epsilon}m_1$ and $Bm_2 = -\frac{1}{2\epsilon}m_2$. The corresponding observable on $S \otimes M$, which will also be denoted B for notational simplicity, takes a product state $|s\rangle|m\rangle \in S \otimes M$ to $|s\rangle B|m\rangle$.

For s a state of S for which we want to weakly measure $\langle s|A|s\rangle$, define a state $f = f(\epsilon)$ of the composite system as the normalization of $(|s\rangle + \epsilon A|s\rangle)|m_1\rangle + (s - \epsilon A|s\rangle)|m_2\rangle$, namely

$$f = f(\epsilon) := \frac{(|s\rangle + \epsilon A|s\rangle)|m_1\rangle + (|s\rangle - \epsilon A|s\rangle)|m_2\rangle}{|(|s\rangle + \epsilon A|s\rangle)|m_1\rangle + (|s\rangle - \epsilon A|s\rangle)|m_2\rangle|}.$$

Then

$$\langle f(\epsilon)|B(\epsilon)|f(\epsilon)\rangle = \langle s|A|s\rangle + O(\epsilon) \quad ,$$

so that $\langle f|B|f\rangle$ approximates $\langle s|A|s\rangle$ to arbitrary accuracy.

To summarize, the average value of $B(\epsilon)$ in the state $f(\epsilon)$ equals (in the limit $\epsilon \rightarrow 0$) the average value of A in the state s . Moreover, the measurements are “weak” in the sense that the state of s after the measurement of B is close to s —the post-measurement state is a mixture of the normalizations of pure states

$s + \epsilon As$ and $s - \epsilon As$, both of which are close to s for small ϵ . The parameter ϵ is a measure of the “weakness” of the measurement: the smaller ϵ is, the less the initial state s is disturbed by the measurement.

This is not the only weak measurement scheme. The seminal paper [1] which introduced weak measurements used a meter space identical to the Hilbert space for a free particle in one dimension, with the “meter” observable B the usual position observable for the free particle¹ so that the “meter” can be visualized as like a pointer moving along a scale. Much of the literature on weak measurements implicitly or explicitly assumes the procedure of [1].

All of the above is uncontroversial, but many aspects of weak measurements *are* controversial. In the literature, weak measurements are often presented in a context in which S is initially in a state s_I (called a *preselected* state) and after the measurement is in a state s_F (called a *postselected* state). One starts with the initial state s_I , performs an “intermediate” measurement of B , and then performs a final measurement to determine whether the state of S is s_F (otherwise it is a mixture of pure states orthogonal to s_F). If the final state is s_F , the result of the intermediate measurement is recorded, otherwise it is discarded. The average of all the recorded measurements of B is called a *weak value* of A .

Since the average of measurements of B *without* postselection equals (in the limits of a large number of measurements with vanishing strength) the average of strong measurements of A , one might be tempted to think that the same would be true *with* postselection. But that overlooks that an intermediate strong measurement of A will project the state of S to an eigenstate of A , which will affect the postselection. Indeed, calculation reveals that the postselected average of strong measurements of A can be no larger than the largest eigenvalue of A (and no smaller than the smallest). Other detailed calculations reveal that the average of measurements of B can turn out to be larger than the largest eigenvalue of A . So, it is controversial whether the postselected average of B may reasonably be considered as a meaningful substitute for any sort of postselected average of A .

Another problem is that different weak measurement procedures can give different postselected meter averages (i.e., averages of B , called “weak values” of A); in fact, in many situations one can obtain arbitrary weak values for A [7] [6]. The seminal paper [1] obtained the weak value

$$\frac{\langle s_F | A | s_I \rangle}{\langle s_F | s_I \rangle} . \quad (1)$$

However, the expression (1) need not be real, and many subsequent authors replace this with its real part.

The motivation of the weak value (1) in [1] is mathematically complicated and its mathematics and physics has been questioned by several authors [3] [4] [2] [5] on different grounds. Moreover, the critics do not always agree with each

¹Actually, they used the momentum observable for the free particle as their meter observable, but Fourier transformation transforms their picture into the one we describe.

other, which is a further indication of the controversial nature of the weak value concept. Many authors use (1) or its real part as if it were the only possible weak value.

References

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